

# Quantum Interferometry and Gravity with Positronium



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*Istituto Nazionale Fisica Nucleare – Milano*  
*On behalf of the QUPLAS Collaboration*

Q U P L A S

QUantum Interferometry, decoherence and gravitational studies with  
Positrons and LASers

- Outline of talk:**
- Basic quantum model of diffraction
  - Fraunhofer and Talbot regimes
  - Incoherence effects
  - Other effects and Decoherence
  - Positrons and Positronium experiments

Home of the Experiment:  
L-NESS Laboratory of the  
Milano Politecnico in Como

# Ps : the truly elementary atom

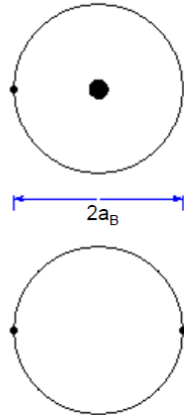
## Energy levels of hydrogen and positronium

$$E_n(H) = -\frac{\mu e^4}{2\hbar n^2} = -\frac{1}{n^2} \times 13.6 \text{ eV}$$

$$\mu_H = \frac{m_e M}{m_e + M} \approx m_e$$

$$\mu_{Ps} = \frac{m_e^2}{2m_e} = \frac{m_e}{2}$$

$$E_n(Ps) = -\frac{1}{n^2} \times 6.8 \text{ eV}$$



A pure QED system where spin-orbit and hyperfine effects are of the same order

The metastable electron-positron bound state can exist in different configurations depending on the relative spin states of the positron and the electron. These are known as para-positronium (p-Ps), with total spin  $S = 0$  and ortho-positronium (o-Ps) with  $S = 1$ .

These spin states have very different lifetimes:

$$|S, m\rangle = |0, 0\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

$$\tau_{p\text{-Ps}} = 125 \text{ ps}$$

$$|S, m\rangle = |1, 0\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$$

$$\tau_{o\text{-Ps}} = 142 \text{ ns}$$

$$|S, m\rangle = |1, 1\rangle = |\uparrow\uparrow\rangle$$

$$|S, m\rangle = |1, -1\rangle = |\downarrow\downarrow\rangle$$

Any process that converts o-Ps to p-Ps is easy to see in lifetime spectra

## 1951: First production of positronium by Martin Deutsch

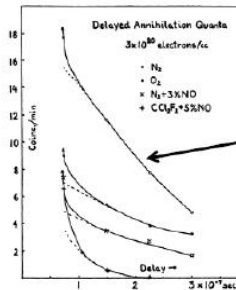
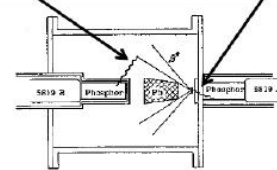


Fig. 1. Decay curves of positrons in several gases. The dotted lines are corrected for time resolution of the instrument.



From M. Deutsch  
Phys. Rev. **82**, 455 (1951)

### Separation of Ps from Radioactive source



## Our systems of interest :

- Electron (an elementary fermion)
- Positron (the antifermion)
- Positronium (Ps, a particle/antiparticle symmetric system)

# The QUPLAS Collaboration

## Università degli Studi di Milano and Infn Milano

S. Castelli, S. Cialdi, M. Giammarchi\*, M. Longhi, G. Maero, Z. Mazzotta,  
S. Olivares, M. Paris, M. Potenza, M. Romè, S. Sala, S. Siccardi, D. Trezzi

## Politecnico Como (Milano)

S. Aghion, M. Bollani (IFN del CNR), G. Consolati, C. Evans, M. Leone, R. Ferragut

## Albert Einstein Center – Laboratory for HEP – University of Bern

A. Ariga, T. Ariga, A. Ereditato, C. Pistillo, P. Scampoli

## Dep.t of Chemistry, University of Bath

K. Edler

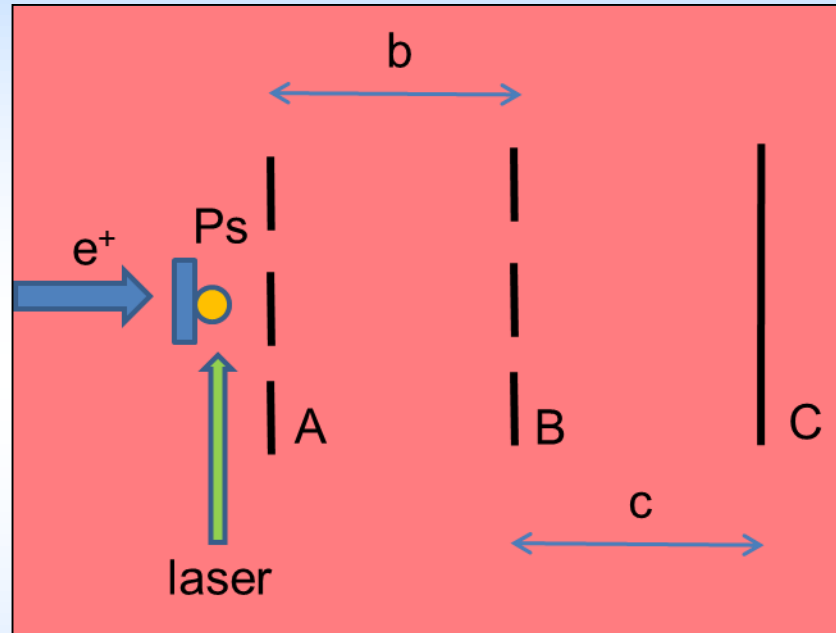
R. Greaves (Los Angeles, formerly at First Point Scientific)

# Introduction to the concept of Quantum Interferometry of Ps

## The typical structure of a Quantum Mechanical Experiment

### Preparation :

- $e^+$  beam
- Ps beam
- Target
- Laser (excitation)
- First grating



### Detection :

- Recording interference pattern
- Projection on measurement eigenstates

Preparation

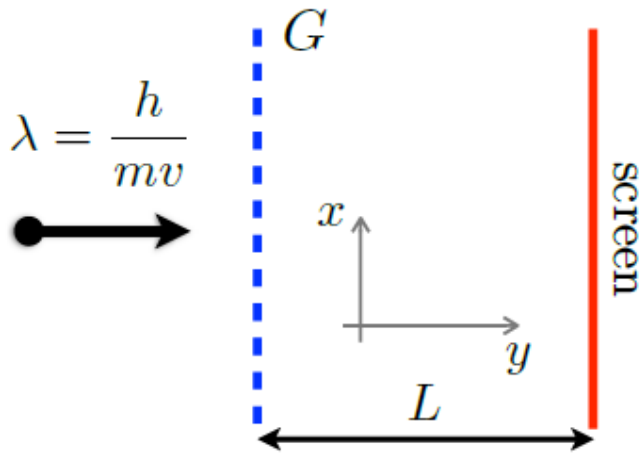
Detection

Non – ideality :  
Incoherence

Interaction  
Propagation  
Interference

Non – ideality :  
Decoherence

# Basic quantum model of diffraction



- **Basic quantum model of diffraction**
- Fraunhofer and Talbot regimes
- Incoherence effects
- Other effects and Decoherence
- Positrons and Positronium experiments

Period D      Slit width a  
Open fraction a/D

De Broglie wave impinging on a grating

- Classical propagation in the y direction
- Schroedinger dynamics in the x direction
- Neglect z-axis diffraction

Interference on the screen given by «Fresnel» integral :

$$I(x) = \left| \psi^{(N)}(x, t = L/v) \right|^2$$

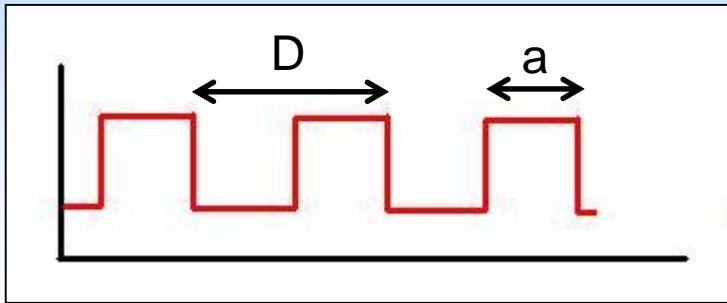
$$\psi^{(N)}(x, t) = \frac{1}{\sqrt{\lambda L}} \int_{-\infty}^{+\infty} \exp \left[ i \frac{\pi}{\lambda L} (x - x')^2 \right] \psi^{(N)}(x', 0) dx'$$

$$\psi^{(N)}(x, t = 0) \approx \sum_{n=1}^N \psi_n(x, t = 0)$$

$$H_{\text{eff}} = \frac{p_x^2}{2m}$$

Possible choices of the initial wavefunction :

Characteristic functions :



$$\psi_n(x,0) = \frac{1}{a} \chi_{\left[\frac{-a}{2} + nD, \frac{+a}{2} + nD\right]}(x)$$

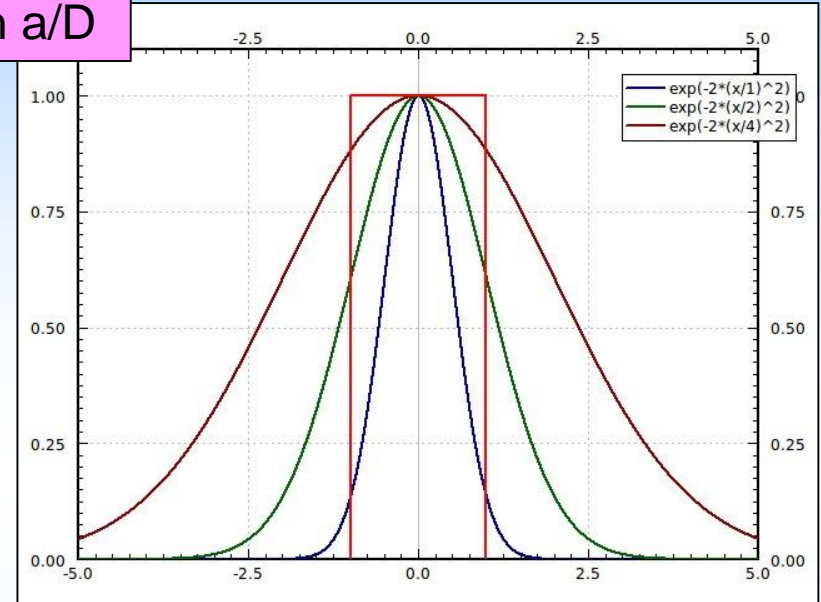
$$\chi_{\Omega}(x) = 1 \text{ if } x \in \Omega$$

$$\chi_{\Omega}(x) = 0 \text{ if } x \notin \Omega$$

Meaning of parameters : crystal clear  
 Total optical analogy  
 Cumbersome calculations

Period D  
 Slit width a  
 Open fraction a/D

Gaussian functions :



$$\psi_n(x,0) = \exp\left[-\frac{(x - nD)^2}{4\sigma^2}\right]$$

Meaning of parameters : from optical analogy

$$\sigma = a / (2\sqrt{2\pi})$$

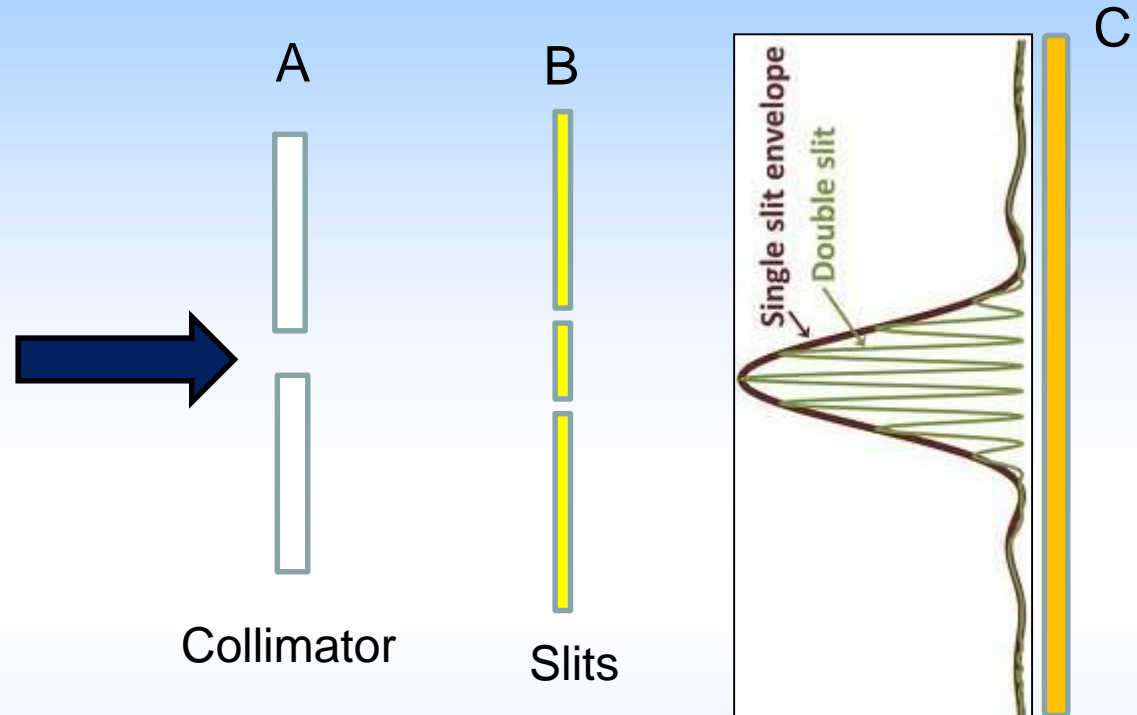
Calculations are easier

## On the choice of the initial single slit wavefunction

Assuming a typical Fraunhofer configuration:

- Radiation beam
- Collimator
- Coherent illumination
- 2 slits
- Detection

One can demonstrate that – similar to classical optics :

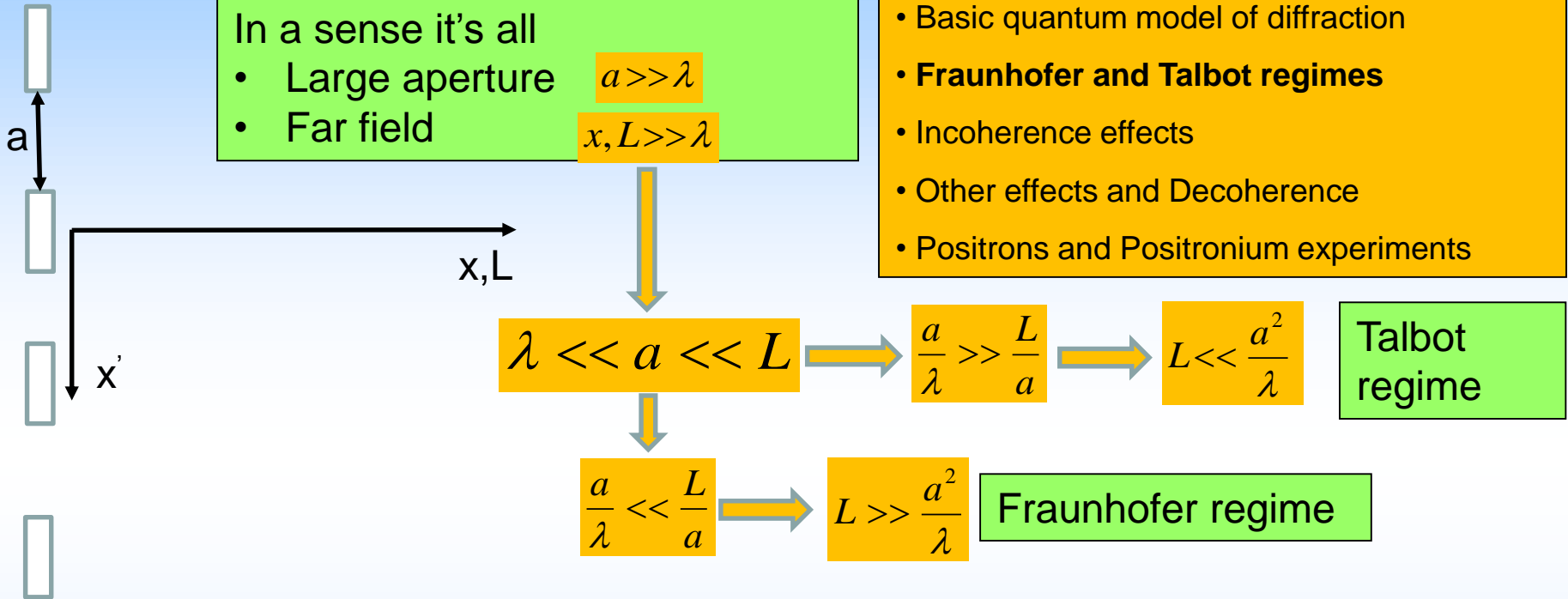


The choice of the single-slit wavefunction impacts only on the envelope of the intensity pattern and not its oscillatory behavior

For more details :

- S. Sala et al., arxiv:1505.01639 [quant-ph]
- S. Sala Master Thesis – University of Milano 2015 – [simone.sala@mi.infn.it](mailto:simone.sala@mi.infn.it)

# Fraunhofer and Talbot regimes



- Basic quantum model of diffraction
- **Fraunhofer and Talbot regimes**
- Incoherence effects
- Other effects and Decoherence
- Positrons and Positronium experiments

Considering the usual expansion, one can define the F parameter

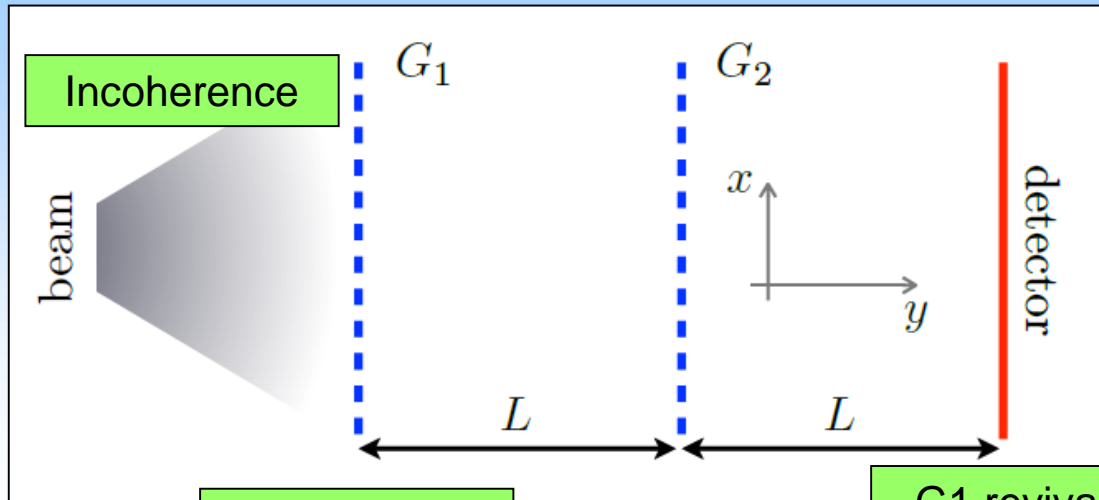
$$F = k \frac{\vec{x}'^2}{\vec{x}^2} = \frac{2\pi}{\lambda} \frac{\vec{x}'^2}{\vec{x}^2} \approx \frac{1}{\lambda} \frac{a^2}{L}$$

$$|\vec{x} - \vec{x}'| = \sqrt{\vec{x}'^2 + \vec{x}^2 - 2\vec{x}\vec{x}'} = |\vec{x}| \sqrt{1 + \frac{\vec{x}'^2}{\vec{x}^2} - 2\frac{\vec{x}\vec{x}'}{\vec{x}^2}}$$

$F \ll 1$  in the Fraunhofer case  
 $F > 1$  in the Talbot case



# The Talbot-Lau Effect (and Talbot «carpets»)



$$L_T = \frac{a^2}{\lambda}$$

Talbot length

The Talbot case

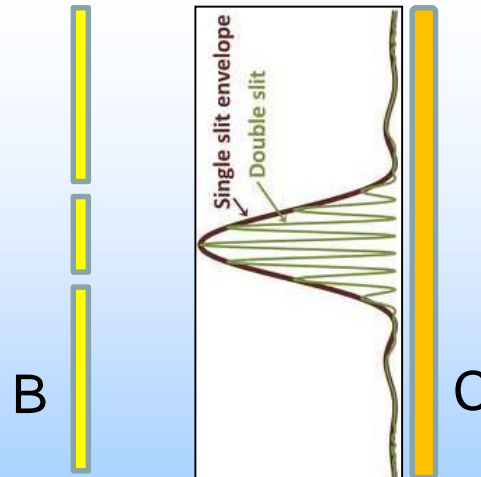
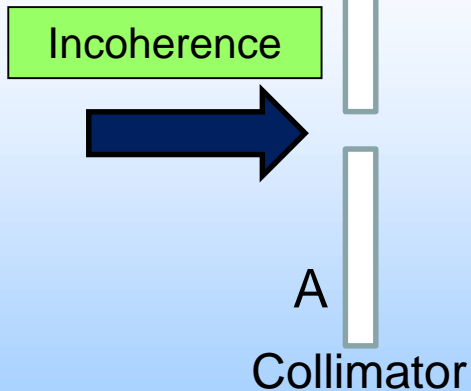
Generates coherence

Interference

G1 revival

$$L = L_T$$

Pattern at C



The Fraunhofer case

$$L \gg L_T$$

$$(BC \gg L_T)$$

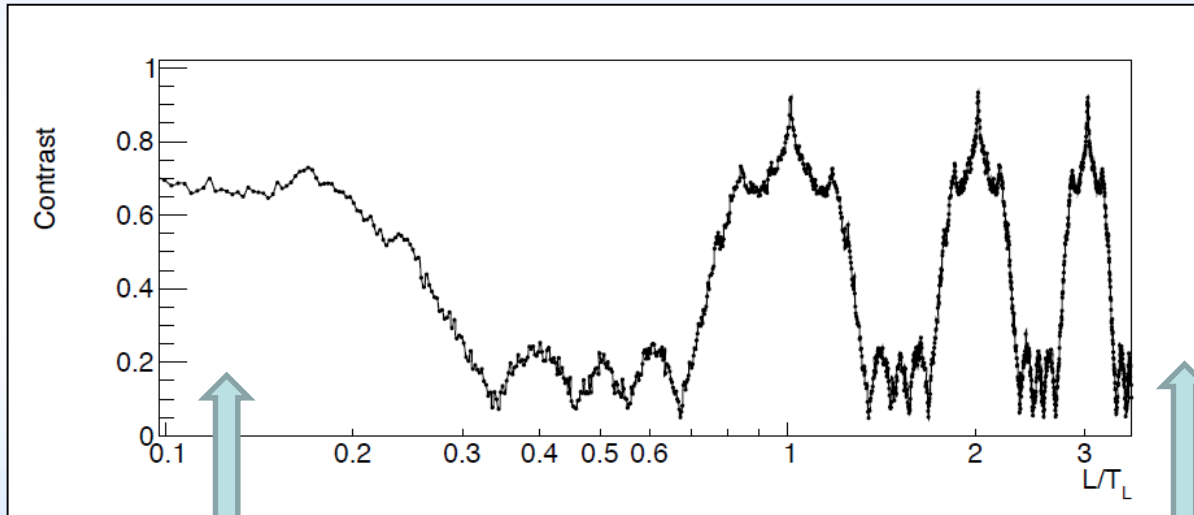
# Talbot carpets

The characteristic pattern of the Talbot effect can be used to make sure the observed effect is the Talbot effect for the specified wavelength

Units  
Talbot  
length

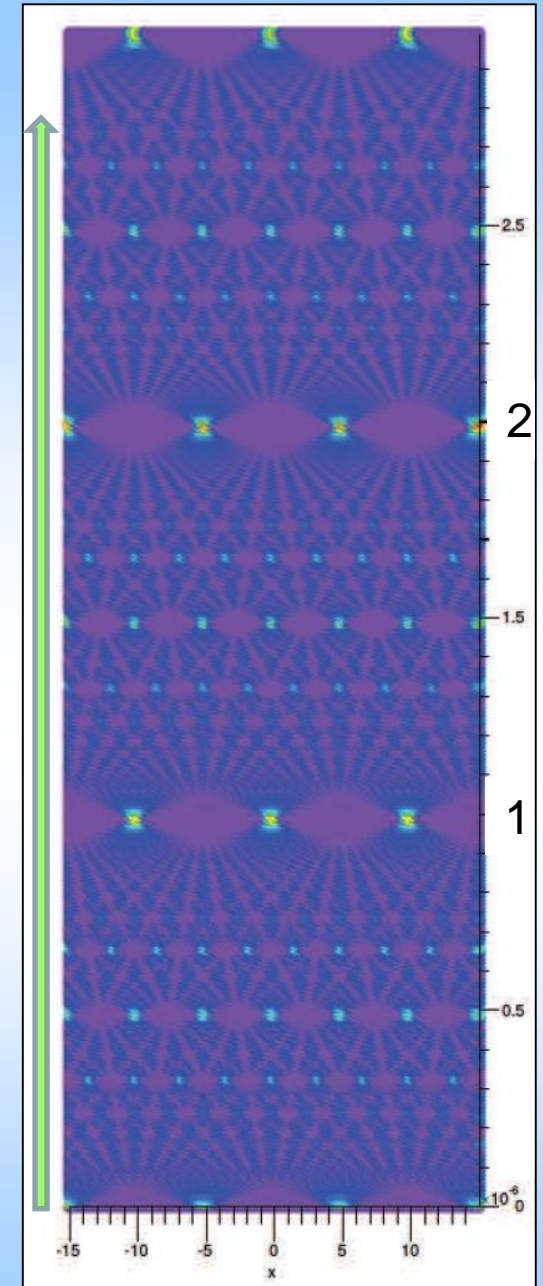
Fringes visibility for the given  
wavelength

$$c = \frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}}$$



«Ballistic» moiré  
regime

Fraunhofer regime  
setting in when  $L \gg L_T$



# Incoherence effects

## We call **Incoherence Effects**...

any effect that could be in principle greatly decreased by a better «classical»

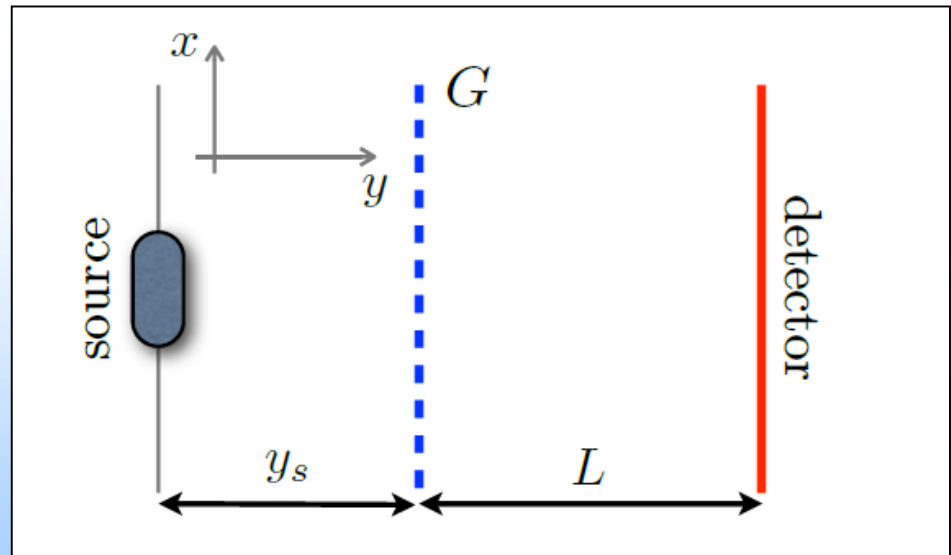
**Preparation** of the Experiment.

- Basic quantum model of diffraction
- Fraunhofer and Talbot regimes
- **Incoherence effects**
- Other effects and Decoherence
- Positrons and Positronium experiments

Incoherence effects are typically due to the source and can often be divided into: transversal (spatial) and longitudinal (time) coherence

Most common examples:

- Spatial extension of the source (typically transverse)
- Non-monochromaticity of the particle velocity spectrum (typically longitudinal)



# Treatment of Incoherence Effects

Physical parameters that can classically fluctuate with a classical distribution. They decrease the visibility

$$\vec{q} = (q_1, q_2, \dots)$$

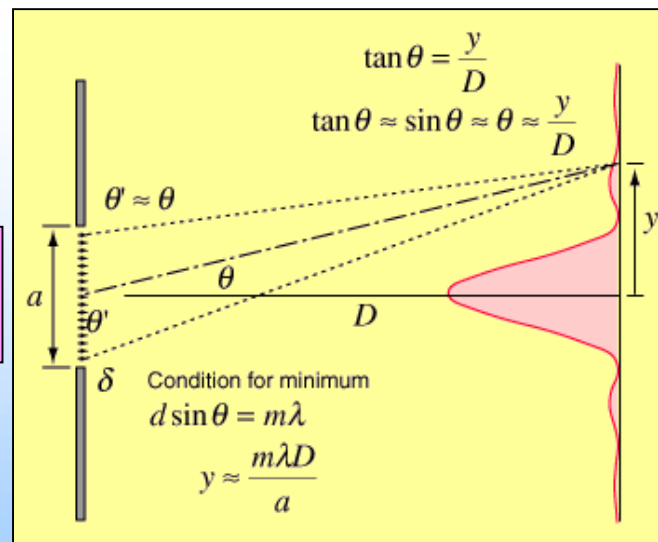
$$p(\vec{q})$$

The effect can be treated by averaging the ideal intensity :

$$I(x, t | \vec{q}) \rightarrow \bar{I}(x, t | \vec{q}) = \int I(x, t | \vec{q}) p(\vec{q}) d\vec{q}$$

The effect sets in, for instance, in limiting the number of actual slits taking coherently part to the interference process

A collimator acting on a beam



An area being coherently illuminated

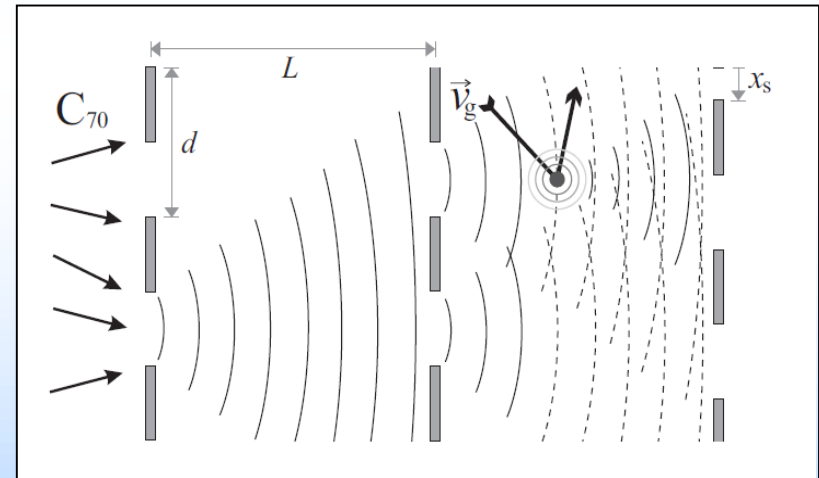
# Other effects and Decoherence

A variety of effects can disturb an interference pattern that cannot be really considered «Incoherence» :

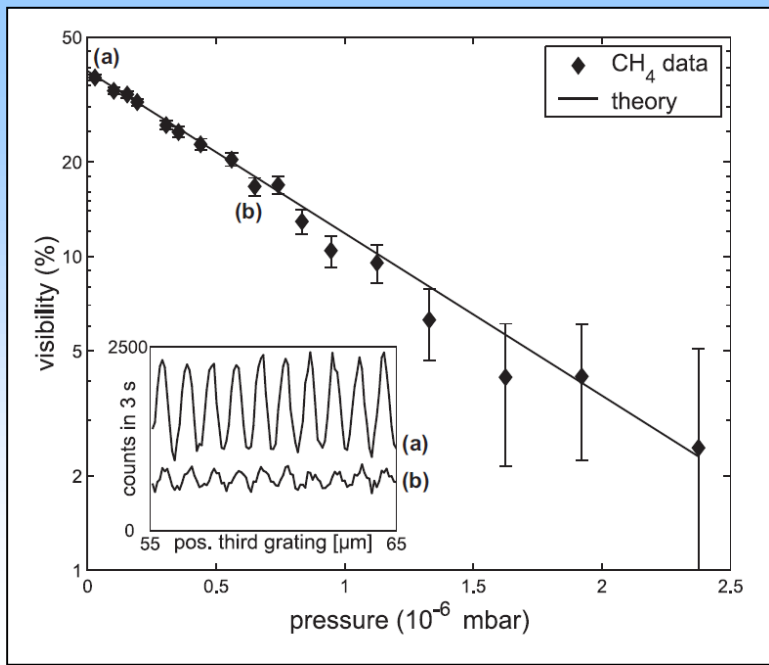
- Particle Decay in Flight
- The physical effect of the grating slits
- Decoherence

Decoherence : loss of the quantum phases between the components of a system in a quantum mechanical superposition. It leads to classical or probabilistically additive behavior. Decoherence occurs when a system interacts with its environment in an irreversible way.

- Basic quantum model of diffraction
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Gas molecules of the background inducing decoherence on  $C_{70}$



Residual pressure in the vacuum chamber is used as a «decoherence parameter»

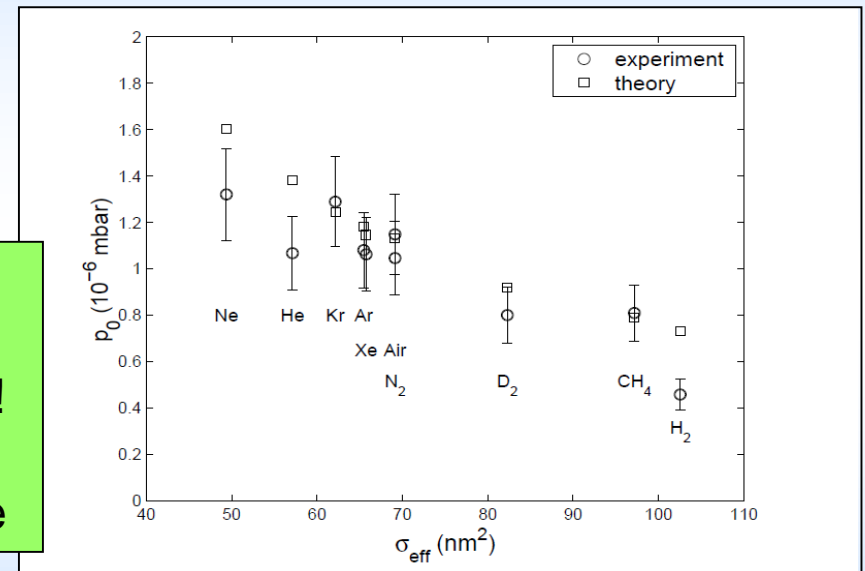
Assumption in the decoherence model :

- Scattering gas massless
- Isotropic velocity distribution of  $\text{C}_{70}$

A Master Equation is solved

Model works for different background gases !

$p$  a good parameter to describe decoherence



## Collisional Decoherence Observed in Matter Wave Interferometry

Klaus Hornberger, Stefan Uttenthaler, Björn Brezger, Lucia Hackermüller, Markus Arndt, and Anton Zeilinger\*

*Universität Wien, Institut für Experimentalphysik, Boltzmannngasse 5, A-1090 Wien, Austria*

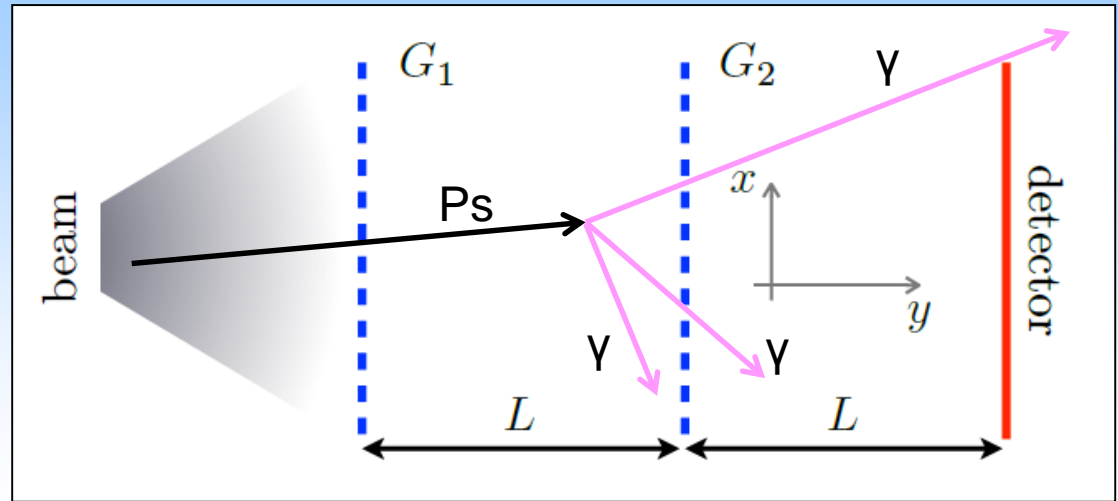
(Dated: March 14, 2003)

# Particle Decay

Positronium states that are useful : ortho-Ps

Decay probability in flight :

$$P(x) \approx \exp(-v\tau/x)$$



The 3-gamma decay of ortho-Ps actually removes the particles from the beam

It does not «blur» the interference pattern in general

So, it is a kind of «attenuation factor» related to a loss of unitarity



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These spin states have very different lifetimes:

$$|S, m\rangle = |0, 0\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

$$\tau_{p-Ps} = 125 \text{ ps}$$

$$|S, m\rangle = |1, 0\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$$

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$$|S, m\rangle = |1, 1\rangle = |\uparrow\uparrow\rangle$$

$$|S, m\rangle = |1, -1\rangle = |\downarrow\downarrow\rangle$$

**Any process that converts o-Ps to p-Ps is easy to see in lifetime spectra**

# Interaction with the slits

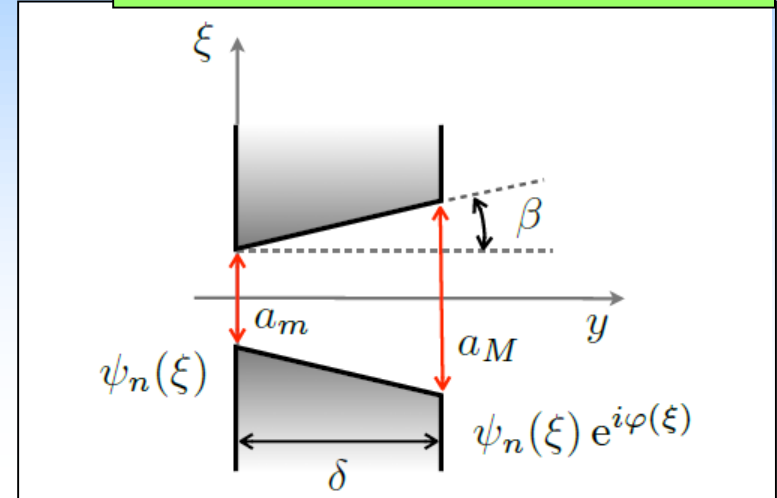
The ideal grating is a perfect 0-1 intensity mask. With no phase effects.

Real gratings have phase effects, instead.

$$\psi_n(\xi, 0) \rightarrow \psi_n(\xi, 0) \exp[i\varphi(\xi)]$$

$$\varphi(\xi) = -\frac{1}{\hbar v} \int V(\xi, y) dy$$

## Beam particle crossing a slit



If the interaction is sufficiently weak, it can be treated by a rescaling of the slit opening  $a \rightarrow a_{\text{eff}}$

Starting from the nominal thickness of 500 nm  
Correction for positrons and antiprotons :

S. Sala et al., arxiv:1505.01639 [quant-ph]

Energy [keV]	$a_{\text{eff}} e^+$ [nm]	$a_{\text{eff}} \bar{p}$ [nm]
0.1	401.3	148.1
1	477.2	285.8
10	497.1	397.4
100	499.7	460.0



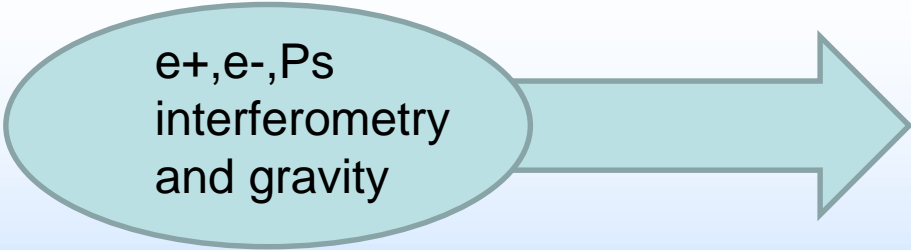
# Positrons and Positronium experiments

Why experiments with positrons and Ps ?

A new type of interferometry

- Positron is a fundamental lepton
- It is the antiparticle of the electron
- Ps is the most fundamental atom

$e^+, e^-, \text{Ps}$   
interferometry  
and gravity



- Basic quantum model of diffraction
- Fraunhofer and Talbot regimes
- Incoherence effects
- Other effects and Decoherence
- **Positrons and Positronium experiments**

**The QUPLAS program**

- QUPLAS-0) Positron Interferometry
- QUPLAS-I) Positronium Interferometry
- QUPLAS-II) Ps gravity

Will describe mostly QUPLAS-0 here

# QUPLAS - 0

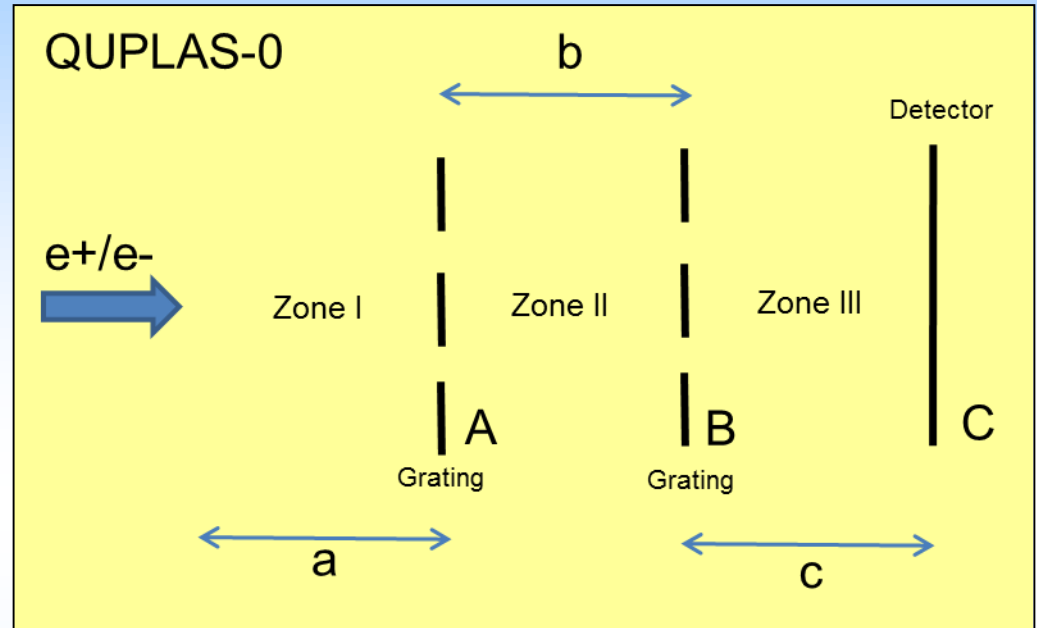
- Interferometry with positrons
- Interferometry with electrons (in the same apparatus)
- Comparison
- A new CPT test

The positron/electron beam :

$T = 10$  keV (typical)

The gratings ( $\sim \mu\text{m}$  thickness):

SiNx based substrates  
Electron Beam Litography



The detector :

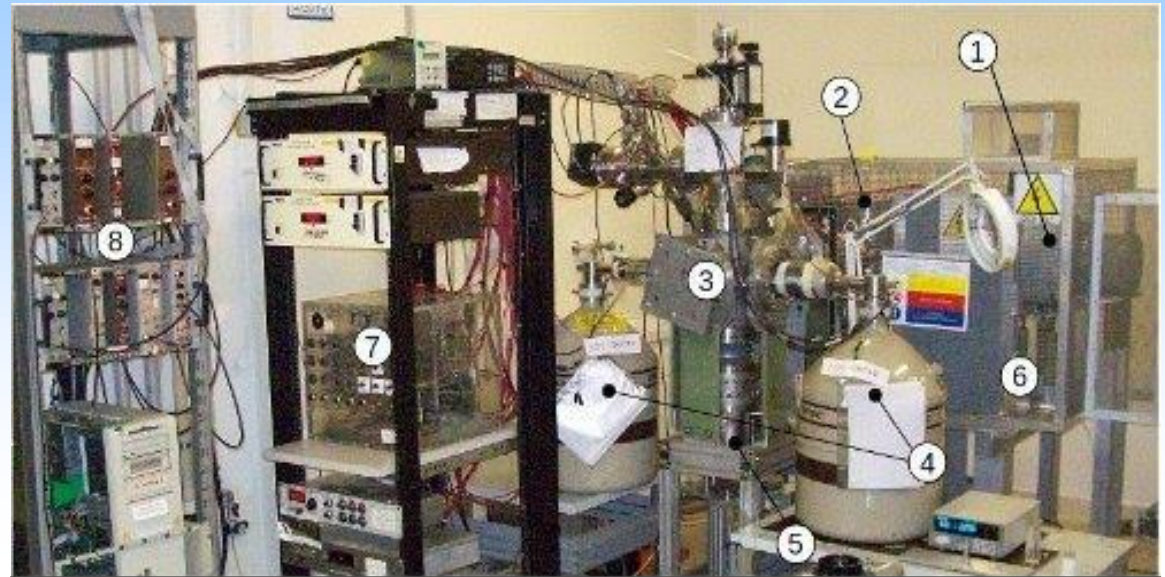
Nuclear Emulsions ( $\sim \mu\text{m}$  resolution)

S. Aghion et al., JINST 8 (2013) P08013.

# The facility: the Como continuous positron beam

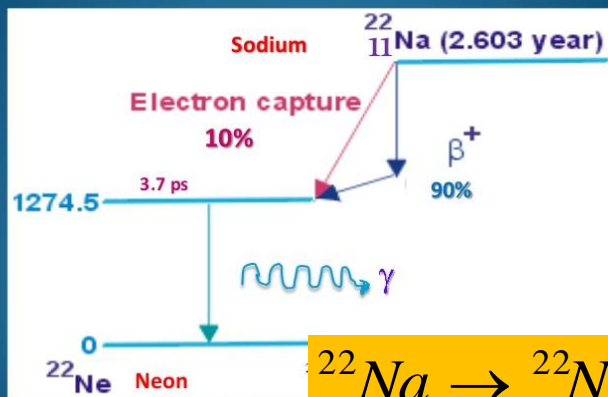
The VPAS Laboratory at the L-Ness Politecnico di Milano at Como Center. (R. Ferragut)

**Slow positron beam.** 1. Radioactive source; 2. Electrostatic optics; 3. Sample chamber; 4. HpGe detectors; 5. Cryostat; 6. High voltage protection cage; 7. Power suppliers; 8. Detector electronics.



<http://www.como.polimi.it/positron>

## Na-22 Decay scheme



Original intensity of the source: 50 mCi  
Current intensity: ~ 13 mCi

Tungsten moderator → reduces the energy from the beta spectrum down to a few eV  
Electrostatic transport → positron beam

Positron beam energy: from a few keV up to 20 keV

Reference value: 10 keV

Intensity:  $\sim 4 \times 10^4$  e<sup>+</sup>/s

$$T = 10 \text{ keV} \quad v = 6 \times 10^7 \text{ m/s}$$

The de Broglie wavelength

$$\lambda = \frac{h}{mv} = 1.2 \times 10^{-11} \text{ m}$$

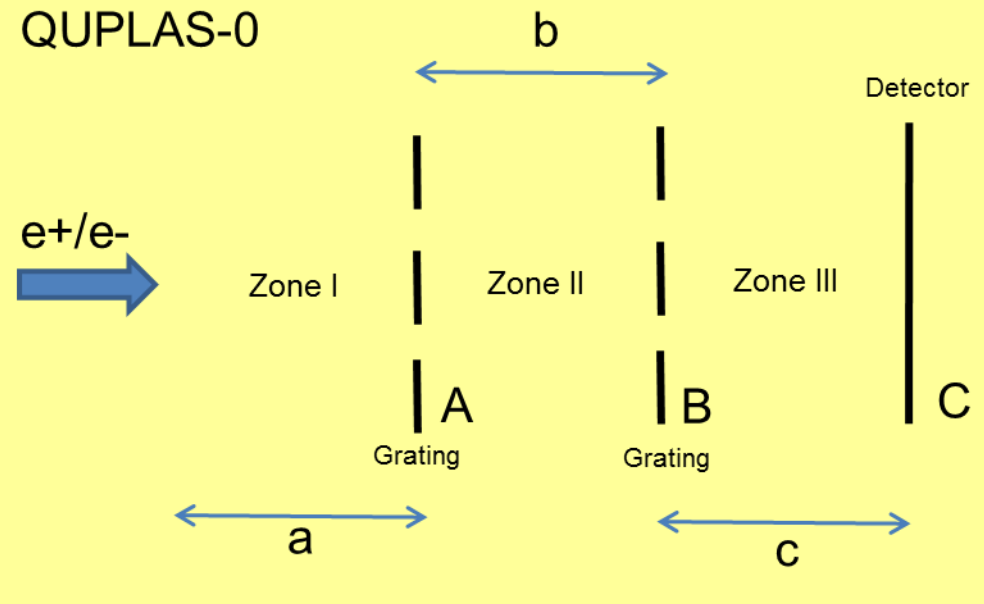
Given a grating with

$$d = 2 \text{ } \mu\text{m}$$

The Talbot length

$$L_T = \frac{d^2}{\lambda} = 33 \text{ cm}$$

QUPLAS-0



One can choose  $b = c = 33$  cm

To have a  $2 \text{ } \mu\text{m}$  periodicity pattern on C

- Setup preparation
- Exposure to the e<sup>+</sup> beam
- Integration on the emulsion detector C

# QUPLAS - I

## Positronium Quantum Interferometry concept

Why?

- Positron Interferometry
- Electron Interferometry
- Positronium Interferometry

An elementary fermion  
The relevant antifermion

The bound fermion-antifermion system  
(also, the simplest atom)

Problems to face :

- Positronium is a neutral atom
- Positronium has a very short lifetime

Detection of the interference pattern is not going to be easy. Ionization required.

Excitation on Rydberg state is necessary.  
Laser excitation required.

# QUPLAS - II

Positronium Gravity : why?

Answer : to test the Weak Equivalence Principle (test of General Relativity)

Universality of Free Fall

## Matter

- Weak Equivalence Principle tested on many different systems
- Torsion Balance Measurement
- $10^{-13}$  level reached

## Antimatter

- g not measured
- Antihydrogen program at CERN (e. g. The AEGIS experiment)
- Aiming at 1% accuracy

## Positronium

- Matter/Antimatter system

?

# Conclusions

Positrons, Electrons, Positronium are interesting !

- Quantum Interference as a key to explore new physics with  $e^+/e^-/Ps$
- QUPLAS is a staged project to tackle these subjects
- QUPLAS-0 : Positron and Electron (charged particles) quantum interference and comparison between them (CPT test)
- QUPLAS-I : Positronium Quantum Interferometry (and a lot of technical development to reach this ambitious goal).
- QUPLAS-II : Positronium Gravity as a test of the Weak Equivalence Principle

Thank you for your attention !

Backup slides



